

# Information Theory – Foundations and Applications

Classical dynamical systems

**Lucas Chibebe Céleri**

Institute of Physics  
Federal University of Goiás

2024 – Basque Center for Applied Mathematics  
University of Basque Country



# Motivation



Definition of **heat** and **work** are problematic in quantum mechanics. **Information** theory can be unambiguously extended to quantum mechanics. The idea of defining thermodynamic quantities in terms of **efficiency** of information processing could pave a new way for a deeper understanding not only of thermodynamics, but also on the limitations for processing information. And, possibly, to a **relativistic** formulation of thermodynamics.

## Goal

To describe a dynamical system as a communication problem and, based on this description, to establish a lower bound on dissipation.

# Some aspects of information and physics<sup>3</sup>



## Jayne's principle

Equilibrium probability distribution maximizes information transfer in the measurement process. Statistical physics can be derived from information theory<sup>1</sup>.

## Landauer's principle

Algorithmic complexity and energy cost of computation (physical implementation of a process) are deeply related. Logical irreversibility implies physical irreversibility<sup>2</sup>.

---

<sup>1</sup>E. T. Jaynes, Phys. Rev. **106**, 620 (1957)

<sup>2</sup>R. Landauer, IBM J. Res. Dev. **5**, 183 (1961).

<sup>3</sup>J. Goold, M. Huber, A. Riera, L. del Rio and P. Skrzypczyk, J. Phys. A **49**, 143001 (2016).

# Some aspects of information and physics



## Phase space

Driven classical Hamiltonian systems<sup>4</sup>.

$$\langle W_{\text{diss}} \rangle \equiv \langle W \rangle - \Delta F = \beta^{-1} D(\rho_t || \tilde{\rho}_t)$$

## Predictive power

Driven and dissipative systems. Unwarranted retention of past information is fundamentally equivalent to energetic inefficiency<sup>5</sup>.

$$\langle W_{\text{diss}}(x_t \rightarrow x_{t+1}) \rangle = \beta^{-1} [I_{\text{mem}}(t) - I_{\text{pre}}(t)]$$

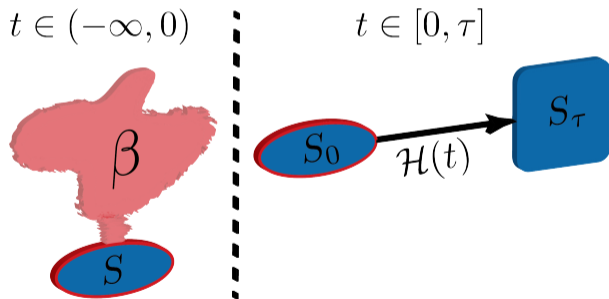
<sup>4</sup>R. Kawai, J. M. R. Parrondo and C. Van den Broeck, Phys. Rev. Lett. **98**, 080602 (2007).

<sup>5</sup>S. Still, D. A. Sivak, A. J. Bell and G. E. Crooks, Phys. Rev. Lett. **109**, 120604 (2012).

## Our problem<sup>6</sup>



Is there any connection between randomness and dissipation? Yes! Based on a description of a dynamical system in terms of communication theory.



<sup>6</sup>M. Capela, M. Sanz, E. Solano and L. C. Céleri, Phys. Rev. E **98**, 052109 (2018)

# Class of systems under consideration



## Elements of the theory

- $\mathcal{H}(s_t; \lambda)$ : System's Hamiltonian.
- $\lambda$ : Set of external controlled parameters.
- $s_t = (q(t), \theta(t))$ : Set of generalized coordinates and canonical conjugate momenta.
- $\Gamma$ : The *finite-dimensional* phase-space.
- The system is initially in the state

$$\rho_0(s_0, \lambda_0) = \frac{e^{-\beta\mathcal{H}(s_0; \lambda_0)}}{Z(\lambda_0)} \quad Z(\lambda) = \int_{\Gamma} ds \exp\{-\beta\mathcal{H}(s; \lambda)\}$$

# Class of systems under consideration



## Dynamics in phase-space

Is deterministic and governed by Hamilton equations

$$\dot{q}_i = \frac{\partial \mathcal{H}(s; \lambda)}{\partial \theta_i} \quad \dot{\theta}_i = -\frac{\partial \mathcal{H}(s; \lambda)}{\partial q_i}$$

- Dynamical system:  $(\Gamma, p, \phi^t)$ .  $(\Gamma, p)$  is a probability space.
- $p : \Sigma \rightarrow [0, 1]$  is the initial probability measure over the sigma-algebra  $\Sigma$ .
- Hamiltonian flow:  $s_t = \phi^t(s_0)$ .  $\phi : \Gamma \rightarrow \Gamma$ .
- The Shannon differential entropy (on the support of the probability density  $\rho_t$ ) is defined as

$$S[\rho_t] = - \int_{\Gamma} ds \rho_t(s) \ln \rho_t(s)$$

# Communication and dynamical systems



The source emits some **discrete symbols** (from some alphabet) to a receiver accordingly with a given probability distribution. The KSE quantifies how random such a process is. The goal here is to define this quantity for dynamical systems.

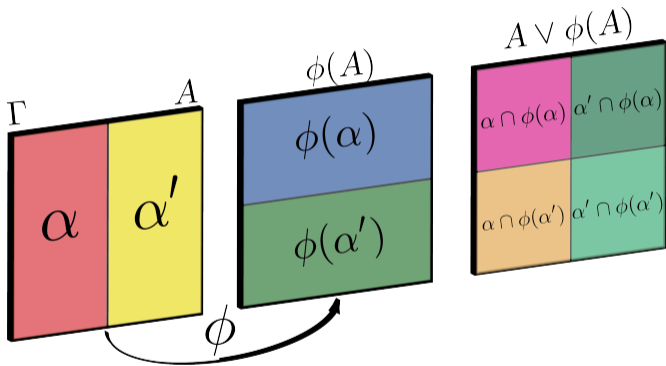
## Partitioning the phase-space

Partition: A collection  $A$  of subsets of the phase-space  $\Gamma$  such that:

- $\forall \alpha, \alpha' \in A, \alpha \cap \alpha' = \emptyset$  if  $\alpha \neq \alpha'$
- $\bigcup_{\alpha \in A} \alpha = \Gamma$



# Refinements of partitions



## Refinement

For partitions  $A$  and  $B$  we define the *refinement*  $A \vee B = \{\alpha \cap \beta \mid \alpha \in A, \beta \in B\}$ .

# The symbols of a dynamical systems



## Discrete time

- Let us consider that time is a discrete variable  $t \in \mathbb{Z}$ .
- Time evolution is generated by iterations of the map  $\phi \equiv \phi^{t=1}$ .
- The phase-space alphabet is constructed as
  - Initial phase-space partition:  $A \rightarrow \phi(A) = \{\phi(\alpha) \mid \alpha \in A\}$
  - The alphabet is provided by the trajectories:  $A, \phi(A), \phi^2(A), \dots$

Kolmogorov-Sinai entropy is then defined for this alphabet.

# The Kolmogorov-Sinai entropy



## Entropy of a partition

$p(\alpha)$ : probability of  $(q, \theta) \in \alpha$ , with  $\alpha \in A$

$$S[A] = - \sum p(\alpha) \ln p(\alpha)$$

## Randomness of the Dynamical System

$$h(\phi) := \sup_{A \in P} \lim_{t \rightarrow \infty} \frac{S \left[ \bigvee_{n=0}^{t-1} \phi^{-n}(A) \right]}{t}, \quad \bigvee_{n=0}^{t-1} \phi^{-n}(A) = A \vee \phi^{-1}(A) \vee \dots \vee \phi^{-t+1}(A)$$

$P$  is the set of all possible finite partitions of  $\Gamma$ .

## Bird's eye view of the proof

- Every initial condition generates a trajectory  $(\alpha_0, \alpha_1, \dots, \alpha_t)$ .
- From this we define the probability for every trajectory and then the conditional *coarse-grained* probability density

$$\rho^{cg}(s_t | \alpha_0, \dots, \alpha_{t-1}) = \sum_{\alpha_t \in A} \frac{p(\alpha_t | \alpha_0, \dots, \alpha_{t-1})}{v(\alpha_t)} \mathbb{1}_{\alpha_t}(s)$$

$\mathbb{1}_{\alpha}(s) = 1$  if  $s \in \alpha$  and  $\mathbb{1}_{\alpha}(s) = 0$  otherwise and  $v(\alpha) = \int_{\alpha} ds$ .

- Compute a lower bound on the phase-space average  $\mathbb{E}S[\rho_t^{cg}]$ .
- Rate in time of  $\mathbb{E}S[\rho_t^{cg}] \rightarrow$  KSE and the rate of  $D(\rho_t || \tilde{\rho}_t) \rightarrow$  rate of  $\langle W_{\text{diss}} \rangle$ .

# Main result



## Lower bound on dissipation rate

$$\beta \overline{\langle W_d \rangle} \geq \beta (\overline{\langle H \rangle} - \overline{F(\lambda_t)}) - \overline{I}_t(A) \quad \overline{I}_t(A) = h(\phi) - \overline{c}_t(A) - \overline{d}_t(A)$$

$$c_t(A) = 1 - \sum_{\alpha_0, \dots, \alpha_t \in A} p(\alpha_t | \alpha_0, \dots, \alpha_{t-1}) \times \tilde{v}(\alpha_{t-1}, \alpha_{t-2}, \dots, \alpha_0 | \alpha_t)$$

$$d_t(A) := - \sum_{\alpha_t \in A} p[\phi^{-t}(\alpha_t)] \ln v[\alpha_t]$$

$F(\lambda_t) := \beta^{-1} \ln Z(\lambda_t)$  is the reference free energy at time  $t$  and the tilde represents backwards quantities.

# Main result: Significance



$$\beta \overline{\langle W_d \rangle} \geq \beta (\overline{\langle H \rangle} - \overline{F(\lambda_t)}) - \overline{I}_t(A) \quad \overline{I}_t(A) = h(\phi) - \overline{c}_t(A) - \overline{d}_t(A)$$

- Hidden information: Difference between the Shannon entropies before and after imposing the coarse-graining.
- $d_0(A)$  is the minimum hidden information:  $S[p(\alpha)] - S[\rho_0] \geq d_0(A) \Rightarrow S[p(\alpha)] - d_0(A)$  is maximum information that is **not** hidden.
- $\overline{\mathbb{E}_{p(\alpha_0, \dots, \alpha_{t-1})} S[p(\alpha_t | \alpha_0, \dots, \alpha_{t-1})]} - \overline{S[\rho_t]} \equiv \overline{I}_t^h$  is the average hidden information.
- $\overline{c}_t(A) + \overline{d}_t(A)$  is the minimum average hidden information:  
 $\overline{I}_t^h \geq \overline{c}_t(A) + \overline{d}_t(A)$ .
- $\overline{I}_t(A)$  is the maximum average information that is **not** hidden ( $A$  is the generating partition): **Information generated by the dynamics.**

# Take home messages



- New tools for studying the thermodynamics of out-of-equilibrium systems based on the understanding of dynamical systems in terms of communication theory.
- In summary, we build a connection between a dynamical quantity, KSE, and a macroscopic physical one, the dissipated work.
- Extension of our results to open systems? Non-Markovianity?<sup>7</sup>
- How about the quantum case? Extension of KSE for quantum stochastic processes<sup>8</sup> and its connections with quantum communication theory should be possible.

---

<sup>7</sup>M. Campisi, P. Hänggi and P. Talkner, Rev. Mod. Phys. **83**, 771 (2011).

<sup>8</sup>G. Lindblad, Commun. Math. Phys. **65**, 281 (1979);

Pollock, Rodríguez-Rosario, Frauenheim, Paternostro and Modi, PRA **97**, 012127 (2018).

# Thank you for your attention

[lucas@qpequi.com](mailto:lucas@qpequi.com)

[www.qpequi.com](http://www.qpequi.com)

